

Noise in Nanometer-scale Electronic Devices

Katsuhiko Nishiguchi and Akira Fujiwara

Abstract

Downsizing electronic devices in integrated circuits increases the noise-related degradation of circuit performance, and thus, it is becoming more important to analyze noise with single-electron resolution. We discuss here the use of a nanometer-scale transistor and capacitor in analyzing thermal noise, one of the most fundamental types of noise in electronic devices, with single-electron resolution. When the capacitor size is substantially reduced, the well-known model of thermal noise is no longer valid, and voltage noise is squeezed.

Keywords: noise, transistor, single electron

1. Introduction

Integrated circuits comprising silicon transistors and other electronic devices are extremely important components in our information technology society. The transistor, one of the most well-known electrical devices, controls the flow of electricity, that is, electrons, for data processing in the integrated circuits. The performance of such integrated circuits has advanced with the downsizing of transistors; for example, the 14-nm generation of transistors will be reached in 2015. However, downsizing also increases the noise-related degradation of circuit performance, and therefore, a microscopic analysis of noise and the development of countermeasures to such noise have become more important than ever before. Indeed, random telegraph noise originating from single-electron trapping in a transistor becomes prominent in small transistors and gives rise to errors in memory circuits. Shot noise originating from the discrete nature of an electron is also expected to become a serious issue in high-speed integrated circuits. Therefore, the downscaling of the transistors increases not only their performance but also the importance of understanding noise microscopically, ultimately with single-electron resolution, in small transistors as well as in other electronic devices.

In this article, we discuss the analysis of thermal

noise, one of the most fundamental types of noise in electronic devices, with single-electron resolution using dynamic random access memory (DRAM) composed of a nanometer-scale transistor and capacitor [1]. Because the analysis was done by using a DRAM, one of the most common devices, we believe that our results are relevant to all electronic devices.

2. Techniques for analyzing thermal noise

One of the simplest and most well-known techniques for analyzing thermal noise is to measure voltage noise in a circuit composed of one resistor and one capacitor, as shown in **Fig. 1**. In this circuit, a massive number of electrons enter and exit the capacitor through the resistor in a random manner with the assistance of thermal energy. This random motion of electrons is thermal noise, and it exists absolutely at a finite temperature. When voltage noise at a node between the resistor and capacitor is measured with an oscilloscope or voltage meter, we can obtain certain information. For example, we can get a histogram of voltage-noise amplitude based on a Gaussian distribution, whose average is voltage V applied to the resistor. We can also get the variance V_{var}^2 of the distribution given by $k_B T/C$, where k_B is the Boltzmann's constant, T is temperature, and C is the capacitance of the capacitor. The fact that V_{var}^2 is proportional to T

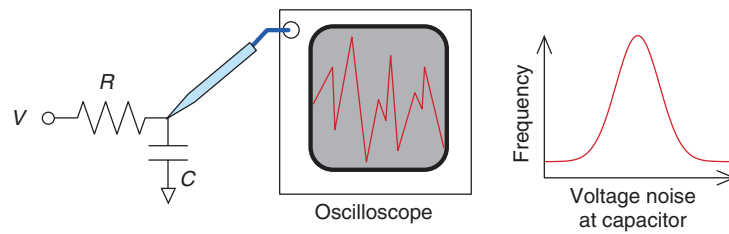


Fig. 1. Conventional technique to measure thermal noise.

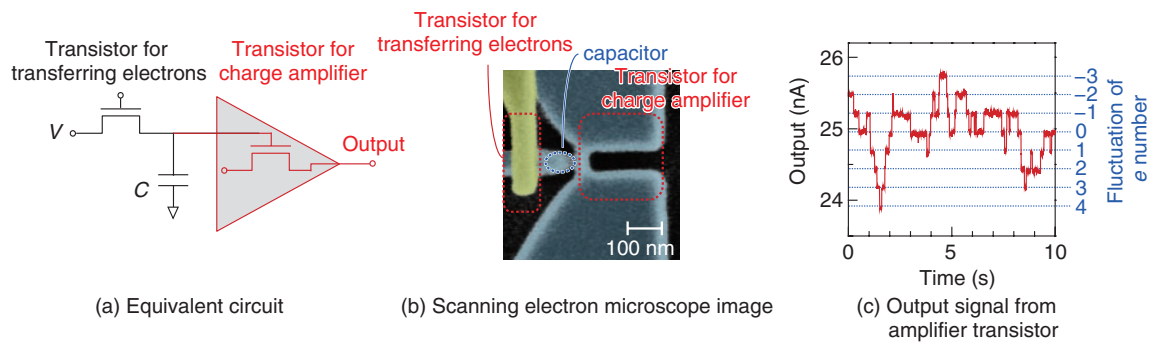


Fig. 2. Measurements of thermal noise with single-electron resolution.

means that the random motion of electrons originates from thermal energy. In contrast, V_{var}^2 increases as C decreases, regardless of the resistor. Therefore, the downsizing of electronic devices in circuits such as memory and analog circuits reduces C and thus increases noise, which leads to an undesirable hindrance of device downsizing. These features can be explained by the well-known Johnson-Nyquist model proposed in 1928 [2, 3].

In our technique for analyzing thermal noise with single-electron resolution, we use instead of a resistor a nanometer-scale transistor for transferring the electrons to the capacitor, as shown in **Fig. 2(a)**, because electrical control of the resistance of the transistor helps us to monitor the motion of individual electrons precisely in real time. In a common circuit, time intervals for the electrons to enter and exit the capacitor are too short for the electrons to be monitored: current I of 1 mA corresponds to the time interval of 1.6×10^{-16} ($=e/I$) seconds, which cannot be measured by any measurement system. However, increasing the resistance of the transistor when it transfers electrons prolongs the interval so that the electron motion can be monitored in real time. These mechanisms and the structure are the same as in a DRAM.

Since the charge of an electron is too small (1.6×10^{-19} C) to be detected by any charge sensor, the tiny signal is amplified by using another transistor [4]: The transistor for the charge amplifier has an extremely small channel (~ 10 nm) and is integrated with the capacitor, as shown in **Fig. 2(b)**. Consequently, the electrons in the capacitor can be detected precisely with single-electron resolution even at room temperature. This success is supported by well-established fabrication techniques for silicon transistors.

Using these features, we can monitor the random motion of individual electrons entering and exiting the capacitor through the transistor in real time, as shown in **Fig. 2(c)**: An output signal from the charge-amplifier transistor changes among discrete values, and one gap between these discrete values corresponds to the charge from one electron. To slow down the electron motion so that it can be monitored, the resistance of the transistor used for the electron transfer is adjusted to be around 10^{20} Ω . This extremely high resistance, which cannot be achieved in conventional resistors and transistors, is possible thanks to our high-quality nanometer-scale transistor.

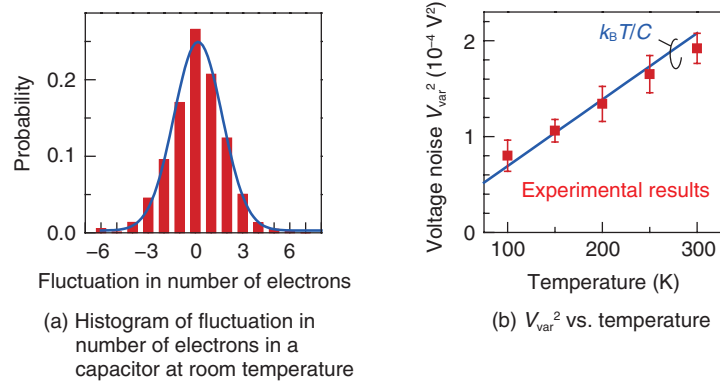


Fig. 3. Thermal noise with single-electron resolution.

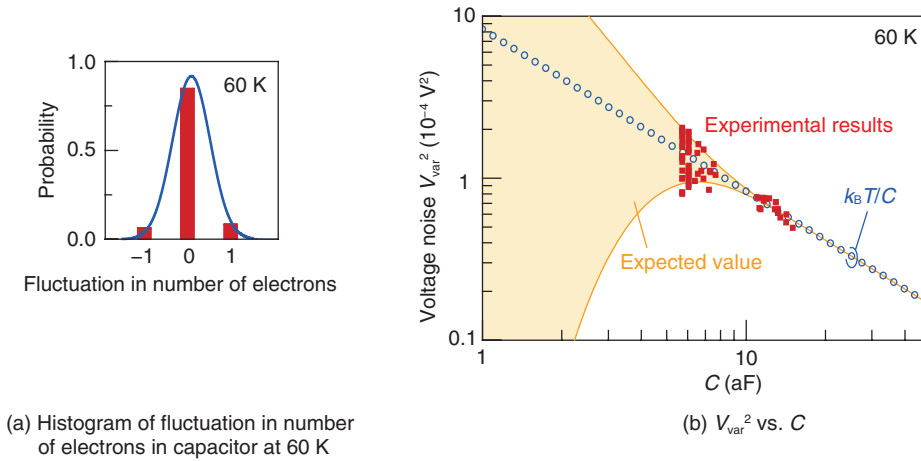


Fig. 4. Thermal noise when charging energy is larger than thermal energy.

3. Thermal noise with single-electron resolution

In analyzing the thermal noise with single-electron resolution, we first discuss the fluctuation, that is, the deviation from the average, of the number of electrons in the capacitor. Because single-electron injection into the capacitor increases the voltage of the capacitor by e/C ($= dV$), the fluctuation of the number of electrons in the capacitor can be converted into a voltage fluctuation, that is, voltage noise, at the capacitor. The dV can be evaluated from the change in the output signal originating from one electron, as shown in Fig. 2(c) [4].

A histogram of the number of electrons in the capacitor is shown in Fig. 3(a). The histogram follows a Gaussian distribution (solid line), as in the case shown in Fig. 1. This histogram gives another

piece of information: dV^2 multiplied by the variance of the distribution shown in Fig. 3(a) corresponds to V_{var}^2 , and this V_{var}^2 follows $k_B T/C$, as shown in Fig. 3(b), which is the same as the case where thermal noise originates from the fluctuation of a massive number of electrons, as explained above and as shown in Fig. 1. Consequently, we can conclude that the well-known Johnson-Nyquist model for thermal noise is adaptable to the random motion of single electrons.

However, when charging energy E_C ($= e^2/2C$) for injecting one electron into the capacitor is larger than thermal energy $k_B T$, the Johnson-Nyquist model is no longer valid. A histogram of the number of electrons in the capacitor at $E_C > k_B T$ is shown in Fig. 4(a). Since the distribution (solid line) becomes sharper, and the available number of electrons becomes

smaller than in the case of $E_C < k_B T$ (Fig. 3(a)), the discrepancy between the experimental results (bars) and expected values (solid line) becomes larger in the form of quantization errors in digital circuits. Indeed, V_{var}^2 ($1.08 \times 10^{-4} \text{ V}^2$) evaluated from the variance of the experimental distribution shown in Fig. 4(a) is smaller than $k_B T/C$ ($1.38 \times 10^{-4} \text{ V}^2$).

This deviation from the Johnson-Nyquist model is similar to the case of electromagnetic radiation from a black body. Energy of the electromagnetic radiation is quantized by multiples of $h\nu$, where h is Planck's constant and ν is the frequency of the electromagnetic radiation. Therefore, when a higher ν makes $h\nu$ larger than thermal energy $k_B T$, the thermal energy cannot assist the emission of electromagnetic waves from the black body. This idea, based on energy quantization proposed by Max Planck, overcomes the ultraviolet catastrophe in the low-frequency region of the spectrum of black-body radiation.

Our results for electrons can also be explained qualitatively by the same model as that for the black-body radiation. Energy for injecting N electrons into the capacitor is given by $(Ne)^2/2C = N^2 E_C$ and quantized because e is a unit charge. Therefore, as in the case of black-body radiation, when E_C is comparable or larger than $k_B T$, the thermal energy cannot assist electrons in entering and exiting the capacitor, which suppresses electron random motion and thus makes V_{var}^2 smaller than the thermal energy.

However, we can observe a unique feature in electron motion unlike in the case of black-body radiation. The average of the number of electrons in the capacitor is given by CV/e , where V is the voltage applied to the transistor as shown in Fig. 2(a). However, for example, when the average number is 0.5, the number of electrons in the capacitor switches between 0 and 1 with the same probability because an electron cannot be divided in half. As a result, the variance of the electron number becomes larger than 0.25, regardless of E_C and $k_B T$. In this sense, the fluctuation in the number of electrons depends on the average number of electrons. When the average number is an integer, the fluctuation becomes minimum; when a fractional part of the average number is 0.5, the fluctuation becomes maximum. The change in voltage noise V_{var}^2 at the capacitor as a function of C is shown in Fig. 4(b). When $E_C < k_B T$ at larger C , V_{var}^2 follows $k_B T/C$ (open circles) and increases as C

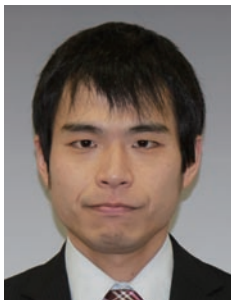
decreases. When $E_C > k_B T$ at smaller C , V_{var}^2 deviates from $k_B T/C$, and the reduction in C increases the discrepancy between V_{var}^2 and $k_B T/C$, as indicated by the shaded area depicting the possible values of V_{var}^2 . It should be noted that the minimum border of the shaded area in Fig. 4(b) decreases with C , behavior that is opposite to that of the conventional $k_B T C$ -limited case depicted by the open circles. In the conventional case valid at $E_C < k_B T$, the reduction in C increases noise, which is a serious issue that hinders the downsizing of electronic devices. However, the case where $E > k_B T$ at much smaller C is preferable to the reduction in C due to noise reduction, which accelerates the downsizing of electronic devices.

4. Conclusion

As the downsizing of electronic devices accelerates, thermal noise will continue to deviate from its well-known behavior. However, this does not mean that any new phenomena will appear. Instead, it means that all of the results observed for single electrons follow well-established thermodynamics, for example, a Boltzmann distribution, and are thus valid for all electronic devices. Therefore, we believe that our results are important in the downsizing of devices to achieve a reduction in the number of electrons in the device. Indeed, using nanometer-scale transistors, we have developed memory and data information circuits that use one electron as one bit of information. In these circuits, errors caused by thermal noise represent one of the most serious issues. However, a wider view is that the minimum energy consumed for computation by data information circuits is governed by the thermal energy. Therefore, the analysis of thermal noise is very important for realizing low-power-consumption circuits.

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Katsuhiko Nishiguchi

Senior Research Scientist, Distinguished Researcher, Nanodevices Group, Physical Science Laboratory, NTT Basic Research Laboratories.

He received his B.E., M.E., and Ph.D. in electrical engineering from Tokyo Institute of Technology in 1998, 2000, and 2002, respectively. Since joining NTT Basic Research Laboratories in 2002, he has been engaged in research on the physics and technology of Si nanometer-scale devices for LSI applications with low power consumption and new functions. He was an invited researcher at the National Center for Scientific Research (CNRS) Lille, France, in September 2008 and also a guest researcher at Delft University of Technology, Delft, the Netherlands in 2012–2013. He received the IUPAP (International Union of Pure and Applied Physics) Young Author Best Paper Award at the International Conference on the Physics of Semiconductors 2000, the Graduate Student Award (Silver) at the Materials Research Society 2000 Fall Meeting, the Young Scientist Award at the Japan Society of Applied Physics (JSAP) Spring Meeting in 2000, the JSAP Outstanding Paper Award 2013, and the Commendation for Science and Technology by the Minister of Education, Culture, Sports, Science and Technology of Japan (Young Scientists' Prize) in 2013. He is a member of the Institute of Electrical and Electronics Engineers (IEEE) and JSAP.



Akira Fujiwara

Senior Distinguished Researcher, Senior Manager of Physical Science Laboratory and Group Leader of Nanodevices Research Group, NTT Basic Research Laboratories.

He received his B.S., M.S., and Ph.D. in applied physics from the University of Tokyo in 1989, 1991, and 1994, respectively. He joined NTT in 1994 and has been engaged in research on silicon nanostructures and their application to nanodevices and single-electron devices. He was a guest researcher at the National Institute of Standards and Technology (NIST), Gaithersburg, MD, USA, in 2003–2004. He was a director of JSAP in 2010–2011 and a visiting professor at Hokkaido University in 2013. He received the International Conference on Solid State Devices and Materials (SSDM) Young Researcher Award in 1998, the SSDM Paper Award in 1999, and JJAP (Japanese Journal of Applied Physics) Paper Awards in 2003, 2006, and 2013. He was awarded the Young Scientist Award from the Ministry of Education, Culture, Sports, Science and Technology in 2006. He was supported by the funding program for Next Generation World-Leading Researchers (NEXT Program), JSPS (Japan Society of Promotion of Science) during 2011–2014. He is a member of JSAP and a senior member of IEEE.
